

Chapter 6. Manipulator Dynamics

11-3-14

Quiz on Nov. 11 on Homework #8

Homework #8. Not collected.

Solve 6.1 (Answer partially given in the textbook), 6.12 (Answer given), 6.16.

Show how (6.32) is derived from (6.15) and (5.45).

Trace the steps taken to derive (6.36) from (6.12).

Verify the formulation of (6.42).

See the Example in Section 6.7 – Two link robot arm with simplifying assumptions.

Check the vector cross multiplications at several places in the solution.

Acceleration of Rigid Body – Definition:

Acceleration of linear velocity vector V_Q in frame {B}

$${}^B\dot{V}_{S_Q} = \frac{d}{dt} {}^B V_Q = \lim_{\Delta t \rightarrow 0} \frac{{}^B V_Q(t + \Delta t) - {}^B V_Q(t)}{\Delta t} \quad (6.1)$$

Acceleration of angular velocity vector ω_Q in frame {B}

$${}^A\dot{\Omega}_Q = \frac{d}{dt} {}^A \Omega_Q = \lim_{\Delta t \rightarrow 0} \frac{{}^A \Omega_Q(t + \Delta t) - {}^A \Omega_Q(t)}{\Delta t} \quad (6.2)$$

Linear Acceleration:

From (5.12),

$${}^A V_Q = \frac{d}{dt} ({}^A R^B Q) = {}^A R^B V_Q + {}^A \Omega_B \times {}^A R^B Q \quad (6.5)$$

Differentiating (6.5) and a term for linear acceleration of the origin of {B},

$${}^A \dot{V}_Q = \frac{d}{dt} ({}^A R^B V_Q) + {}^A \dot{\Omega}_B \times {}^A R^B Q + {}^A \Omega_B \times \frac{d}{dt} ({}^A R^B Q) \quad (6.7)$$

$$= ({}^A R^B \dot{V}_Q + {}^A \Omega_B \times {}^A R^B V_Q) + {}^A \dot{\Omega}_B \times {}^A R^B Q + {}^A \Omega_B \times ({}^A R^B V_Q + {}^A \Omega_B \times {}^A R^B Q) \quad (6.8)$$

With the linear acceleration of {B} Orig

$${}^A \dot{V}_Q = {}^A \dot{V}_{BOrg} + {}^A \Omega_B \times {}^A R^B V_Q + 2 {}^A \Omega_B \times {}^A R^B V_Q + {}^A R^B \dot{V}_Q + {}^A \dot{\Omega}_B \times {}^A R^B Q + {}^A \Omega_B \times ({}^A \Omega_B \times {}^A R^B Q) \quad (6.10)$$

When ${}^B Q$ is constant,

$${}^A \dot{V}_Q = {}^A \dot{V}_{BOrg} + {}^A R^B \dot{V}_Q + {}^A \dot{\Omega}_B \times {}^A R^B Q + {}^A \Omega_B \times ({}^A \Omega_B \times {}^A R^B Q) \quad (6.12)$$

Angular Acceleration:

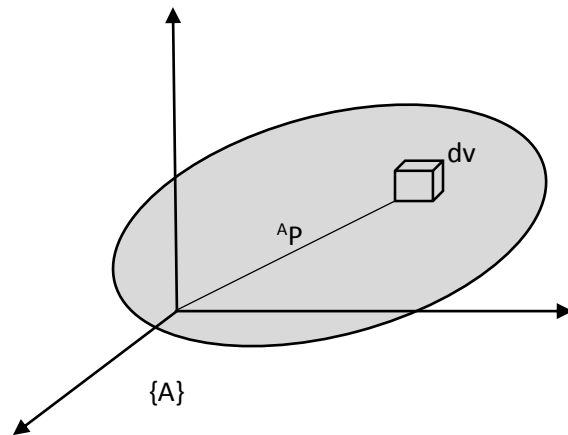
To find the angular acceleration of {C} w.r.t. {A}, differentiate

$${}^A\Omega_C = {}^A\Omega_B + {}^A R^B \Omega_C \quad (6.13)$$

$${}^A\dot{\Omega}_C = {}^A\dot{\Omega}_B + \frac{d}{dt}({}^A R^B \Omega_C) = {}^A\dot{\Omega}_B + {}^A R^B \dot{\Omega}_C + {}^A\Omega_B \times {}^A R^B \Omega_C \quad (6.15)$$

Rigid Body Mass Distribution

Inertia tensor – Describes the distribution of the mass around the center of a rigid body.



${}^A P$ is the location vector of the differential volume dv .

Inertia Tensor of {A}:

$${}^A I = \begin{bmatrix} I_{xx} & -I_{xy} & -I_{xz} \\ -I_{xy} & I_{yy} & -I_{yz} \\ -I_{xz} & -I_{yz} & I_{zz} \end{bmatrix}$$

Mass moment of inertia:

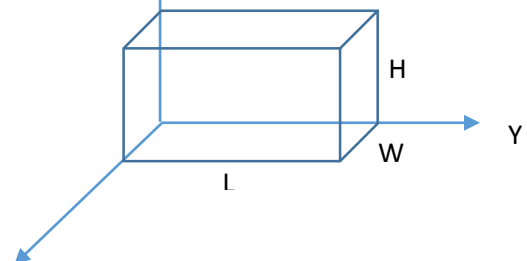
$$I_{xx} = \iiint_V (y^2 + z^2) \rho dv \quad I_{yy} = \iiint_V (x^2 + z^2) \rho dv \quad I_{zz} = \iiint_V (x^2 + y^2) \rho dv$$

$$I_{xy} = \iiint_V xy \rho dv \quad I_{xz} = \iiint_V xz \rho dv \quad I_{yz} = \iiint_V yz \rho dv$$

Example 6.1

$$I_{xx} = \int_0^h \int_0^l \int_0^w (r^2 - x^2) \rho dx dy dz = \int_0^h \int_0^l \int_0^w (y^2 + z^2) \rho dx dy dz = \frac{m}{3} (l^2 + h^2)$$

$$I_{xy} = \int_0^h \int_0^l xy \rho dx dy dz = \frac{m}{4} wl$$



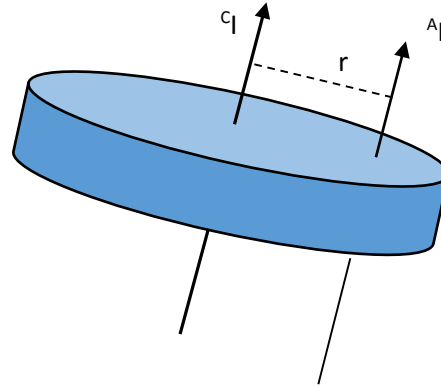
Parallel Axis Theorem:

The moment of inertia at the center of the mass is at the minimum quantity along the axis of rotation. The moment of inertia of any axis parallel to the axis of rotation is given by

$${}^A I_{zz} = {}^C I_{zz} + m \cdot r_c^2$$

where r_c = the distance from the axis in {A} to the center of the mass in {C} and m = the point mass at the center.

Inertial tensor of a mass in frame {A} w.r.t. frame {C} with its origin at the center of the mass.



$${}^A I_{zz} = {}^C I_{zz} + m(x_c^2 + y_c^2) = {}^C I_{zz} + m r_c^2$$

$${}^A I_{xy} = {}^C I_{xy} - m x_c y_c$$

$${}^A P_c = [x_c \quad y_c \quad z_c]^T$$

- Location of the center of mass in {A}.

The frame {A} has its origin at ${}^A P_c = \frac{1}{2} [w \quad l \quad h]^T$

$${}^C I_{zz} = \frac{m}{12} (w^2 + l^2) \quad {}^C I_{xy} = 0$$

$${}^C I = \begin{bmatrix} {}^C I_{xx} & 0 & 0 \\ 0 & {}^C I_{yy} & 0 \\ 0 & 0 & {}^C I_{zz} \end{bmatrix}$$

Example 6.2

Newton's Equation on Force: $F = m \dot{v}_C$ at the center of mass

Euler's Equation on Moment: $N = {}^C I \dot{\omega} + \omega \times {}^C I \omega$ at the center of mass

${}^C I$ = inertia tensor in frame {C} with its origin at the mass center

Newton-Euler Dynamic Equations

Derivation of angular acceleration

Forward angular velocity propagation

$${}^{i+1}\dot{\omega}_{i+1} = {}^{i+1}R^i \dot{\omega}_i + \dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} \quad \text{from (5.45)}$$

$${}^A\dot{\Omega}_C = {}^A\dot{\Omega}_B + {}^A R^B \dot{\Omega}_C + {}^A\Omega_B \times {}^A R^B \Omega_C \quad \text{from (6.15)}$$

$$(6.15')$$

$$(6.15'')$$

Follow the derivation of (6.32) from (6.15) and (5.45)

$${}^{i+1}\dot{\omega}_{i+1} = {}^{i+1}R^i \dot{\omega}_i + {}^{i+1}R^i \omega_i \times \dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} + \ddot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} \quad (6.32)$$

For prismatic joints, ${}^i\omega_i = \dot{\theta}_i = 0$, so

$${}^{i+1}\dot{\omega}_{i+1} = {}^{i+1}R^i \dot{\omega}_i$$

Derivation of linear acceleration

From (6.12) and by taking similar steps to derive the angular acceleration,

$${}^A V_Q = \frac{d}{dt} ({}^A R^B Q) = {}^A R^B V_Q + {}^A\Omega_B \times {}^A R^B Q$$

$${}^A \dot{V}_Q = {}^A \dot{V}_{Borg} + {}^A R^B \dot{V}_Q + {}^A \dot{\Omega}_B \times {}^A R^B Q + {}^A\Omega_B \times ({}^A\Omega_B \times {}^A R^B Q)$$

Setting {A}={i+1} and {B}={i} and factoring out ${}^{i+1}R$,

$${}^{i+1}\dot{v}_{i+1} = {}^{i+1}R^i [\dot{v}_i + \dot{\omega}_i \times {}^i P_{i+1} + {}^i\omega_i \times ({}^i\omega_i \times {}^i P_{i+1})] \quad (6.34)$$

For prismatic joints, add two more terms to (6.34) per (6.10)

$${}^{i+1}\dot{v}_{i+1} = {}^{i+1}R^i [\dot{v}_i + \dot{\omega}_i \times {}^i P_{i+1} + {}^i\omega_i \times ({}^i\omega_i \times {}^i P_{i+1})] + 2{}^{i+1}\omega_{i+1} \times \dot{d}_{i+1} {}^{i+1}\hat{Z}_{i+1} + \ddot{d}_{i+1} {}^{i+1}\hat{Z}_{i+1} \quad (6.35)$$

Linear acceleration of the center of mass, from (6.12)

Trace the steps taken in applying (6.12),

$${}^i\dot{v}_{Ci} = {}^i\dot{\omega}_i \times {}^i P_{Ci} + {}^i\omega_i \times ({}^i\omega_i \times {}^i P_{Ci}) + {}^i\dot{v}_i \quad (6.36)$$

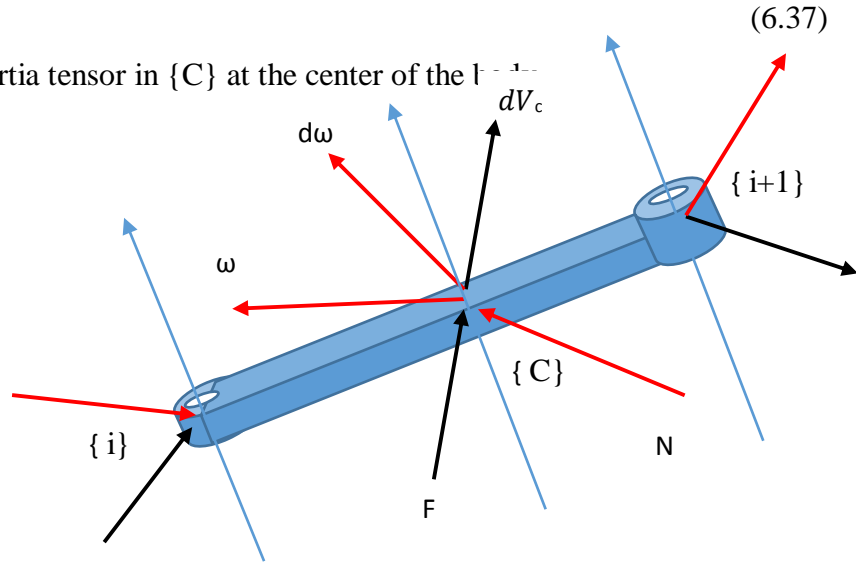
The inertial force and torque acting at the center of the mass:

From (6.32) and (6.36)

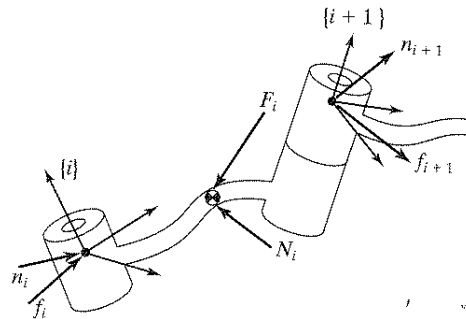
$$F_i = m\dot{v}_{c_i}$$

$$N_i = {}^C I \dot{\omega}_i + \omega_i \times {}^C I \omega_i$$

where 3×3 ${}^C I$ is the inertia tensor in $\{C\}$ at the center of the mass



H7 Backward Iteration for Joint Forces and Torque



Force and torque balance equations at the center of mass of link i :

$${}^i F_i = {}^i f_i - {}^i R^{i+1} f_{i+1} \quad (6.38)$$

$${}^i N_i = {}^i n_i - {}^i n_{i+1} + ({}^i P_i - {}^i P_{C_i}) \times {}^i f_i - ({}^i P_{i+1} - {}^i P_{C_i}) \times {}^i f_{i+1} \quad (6.39)$$

$${}^i P_i = 0$$

Rearranging the equations and adding rotations;

$${}^i f_i = {}^i R^{i+1} f_{i+1} + {}^i F_i \quad (6.41)$$

$${}^i n_i = {}^i N_i + {}^i R^{i+1} n_{i+1} + {}^i P_{C_i} \times {}^i F_i + {}^i P_{i+1} \times {}^i R^{i+1} f_{i+1} \quad (6.42)$$

Finally, the joint torque is the Z component of the vector representing the inertial torque:

$$\tau_i = {}^i n_i^T \hat{Z}_i \quad (6.43)$$

For prismatic joints, using τ to denote force:

$$\tau_i = {}^i f_i^T \hat{Z}_i \quad (6.44)$$

Forward and backward iterations: Eq (6.45)-(6.53)

Forward - Link velocities and accelerations via the Newton-Euler (6.31)-(6.37).

$${}^{i+1}\omega_{i+1} = {}^i R^{i+1} \omega_i + \dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} \quad (6.45)$$

$${}^{i+1}\dot{\omega}_{i+1} = {}^i R^{i+1} \dot{\omega}_i + {}^i R^{i+1} \omega_i \times \dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} + \ddot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} \quad (6.46)$$

$${}^{i+1}\dot{v}_{i+1} = {}^i R^{i+1} [\dot{v}_i + \dot{\omega}_i \times {}^i P_{i+1} + \omega_i \times ({}^i \omega_i \times {}^i P_{i+1})] \quad (6.47)$$

$${}^{i+1}\dot{v}_{C_{i+1}} = \dot{\omega}_i \times {}^{i+1} P_{i+1} + {}^{i+1} \omega_{i+1} \times ({}^{i+1} \omega_{i+1} \times {}^{i+1} P_{i+1}) + {}^{i+1} \dot{v}_{i+1} \quad (6.48)$$

$${}^{i+1} F_{i+1} = m_{i+1} {}^{i+1} \dot{v}_{C_{i+1}} \quad (6.49, 6.50)$$

$${}^{i+1} N_{i+1} = {}^{C_{i+1}} I_{i+1} {}^{i+1} \dot{\omega}_{i+1} +$$

$${}^{i+1} N_{i+1} = {}^{C_{i+1}} I_{i+1} {}^{i+1} \dot{\omega}_{i+1} + {}^{i+1} \omega_{i+1} \times {}^{i+1} I_{i+1} {}^{i+1} \omega_{i+1}$$

Backward - Find joint forces and torques via (6.38)-(6.44).

$${}^i f_i = {}^i I_{i+1} {}^i R^{i+1} f_{i+1} + F_i \quad (6.51)$$

$${}^i n_i = {}^i N_i + {}^i R^{i+1} n_{i+1} + {}^i P_i \times {}^i F_i + {}^i P_{i+1} \times {}^i R^{i+1} f_{i+1} \quad (6.52)$$

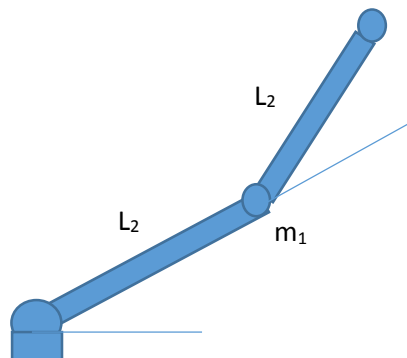
$$\tau_i = {}^i n_i^T {}^i \hat{Z}_i \quad (6.53)$$

See the Example in Section 6.7 – Simplified two link robot arm.

Check the vector cross multiplications at several places in the solution.

$$\mathbf{u} \times \mathbf{v} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ u_1 & u_2 & u_3 \\ v_1 & v_2 & v_3 \end{vmatrix} \quad m_2$$

$$\mathbf{u} \times \mathbf{v} = \begin{vmatrix} u_2 & u_3 \\ v_2 & v_3 \end{vmatrix} \mathbf{i} - \begin{vmatrix} u_1 & u_3 \\ v_1 & v_3 \end{vmatrix} \mathbf{j} + \begin{vmatrix} u_1 & u_2 \\ v_1 & v_2 \end{vmatrix} \mathbf{k}$$



For Link 1:

$${}^1\omega_1 = \omega_0 + \dot{\theta}_1 {}^1\hat{Z}_1 = \begin{bmatrix} 0 & 0 & \dot{\theta}_1 \end{bmatrix}^T$$

$${}^1\dot{\omega}_1 = \ddot{\theta}_1 {}^1\hat{Z}_1 = \begin{bmatrix} 0 & 0 & \ddot{\theta}_1 \end{bmatrix}^T$$

$${}^1\dot{v}_1 = {}^1R^i \dot{v}_i = \begin{bmatrix} c_1 & s_1 & 0 \\ -s_1 & c_1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ g \\ 0 \end{bmatrix} = \begin{bmatrix} gs_1 \\ gc_1 \\ 0 \end{bmatrix}$$

$${}^1\dot{v}_{C_1} = {}^1\dot{\omega}_1 \times {}^1P_{C_1} + {}^1\omega_1 \times ({}^1\omega_1 \times {}^1P_{C_1}) + {}^1\dot{v}_1 = \begin{bmatrix} i & j & k \\ 0 & 0 & \ddot{\theta}_1 \\ l_1 & 0 & 0 \end{bmatrix} + {}^1\omega_1 \times \begin{bmatrix} i & j & k \\ 0 & 0 & \dot{\theta}_1 \\ l_1 & 0 & 0 \end{bmatrix} + \begin{bmatrix} gs_1 \\ gc_1 \\ 0 \end{bmatrix} = \begin{bmatrix} -l_1\dot{\theta}_1^2 + gs_1 \\ l_1\ddot{\theta}_1 + gs_1 \\ 0 \end{bmatrix}$$

$${}^1F_1 = m_1 {}^1\dot{v}_1 = m_1 \begin{bmatrix} -l_1\dot{\theta}_1^2 + gs_1 \\ l_1\ddot{\theta}_1 + gs_1 \\ 0 \end{bmatrix}$$

$${}^1N_1 = {}^{C_1}I_{i+1} {}^1\dot{\omega}_1 + {}^1\omega_1 \times {}^{i+1}I_{i+1} {}^1\omega_1 = 0 \cdot {}^1\dot{\omega}_1 + {}^1\omega_1 \times 0 \cdot {}^1\omega_1 = \begin{bmatrix} 0 & 0 & 0 \end{bmatrix}^T$$

For Link 2:

$${}^2\omega_2 = {}^2R^1 \omega_1 + \dot{\theta}_2 {}^2\hat{Z}_2 = \begin{bmatrix} c_2 & s_2 & 0 \\ -s_2 & c_2 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 + \dot{\theta}_2 \end{bmatrix}$$

$${}^2\dot{\omega}_2 = \begin{bmatrix} 0 \\ 0 \\ \ddot{\theta}_1 + \ddot{\theta}_2 \end{bmatrix}$$

$${}^2\dot{v}_2 = {}^2R^1 [{}^1\dot{v}_1 + {}^1\dot{\omega}_1 \times {}^1P_2 + {}^1\omega_1 \times ({}^1\omega_1 \times {}^1P_2)] = {}^2R^2 \dot{v}_{C_1} = \begin{bmatrix} c_2 & s_2 & 0 \\ -s_2 & c_2 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} -l_1\dot{\theta}_1^2 + gs_1 \\ l_1\ddot{\theta}_1 + gs_1 \\ 0 \end{bmatrix}$$

$${}^2\dot{v}_{C_2} = {}^1\dot{\omega}_1 \times {}^2P_{C_2} + {}^2\omega_2 \times ({}^2\omega_2 \times {}^2P_{C_2}) + {}^2\dot{v}_2 = \begin{bmatrix} i & j & k \\ 0 & 0 & \ddot{\theta}_1 \\ l_2 & 0 & 0 \end{bmatrix} + {}^2\omega_2 \times \begin{bmatrix} i & j & k \\ 0 & 0 & \dot{\theta}_1 + \dot{\theta}_2 \\ l_2 & 0 & 0 \end{bmatrix} + {}^2\dot{v}_2$$

$${}^2F_2 = m_2 {}^2\dot{v}_{C_2}$$

$${}^2N_2 = {}^{C_2}I_2 {}^2\dot{\omega}_2 + {}^2\omega_2 \times {}^{C_2}I_2 {}^2\omega_2 = [0] {}^2\dot{\omega}_2 + {}^2\omega_2 \times [0] {}^2\omega_2 = \begin{bmatrix} 0 & 0 & 0 \end{bmatrix}^T$$

The backward calculations of ${}^2f_2, {}^2n_2, {}^1f_1, {}^1n_1$ for Links 2 and 1 can be carried out similarly.

The joint torques τ_1, τ_2 are the Z components of ${}^1n_1, {}^2n_2$.